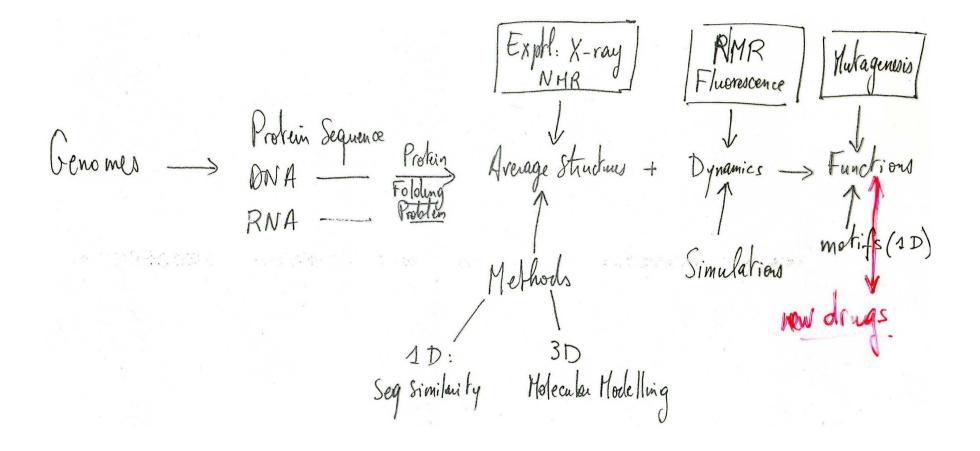
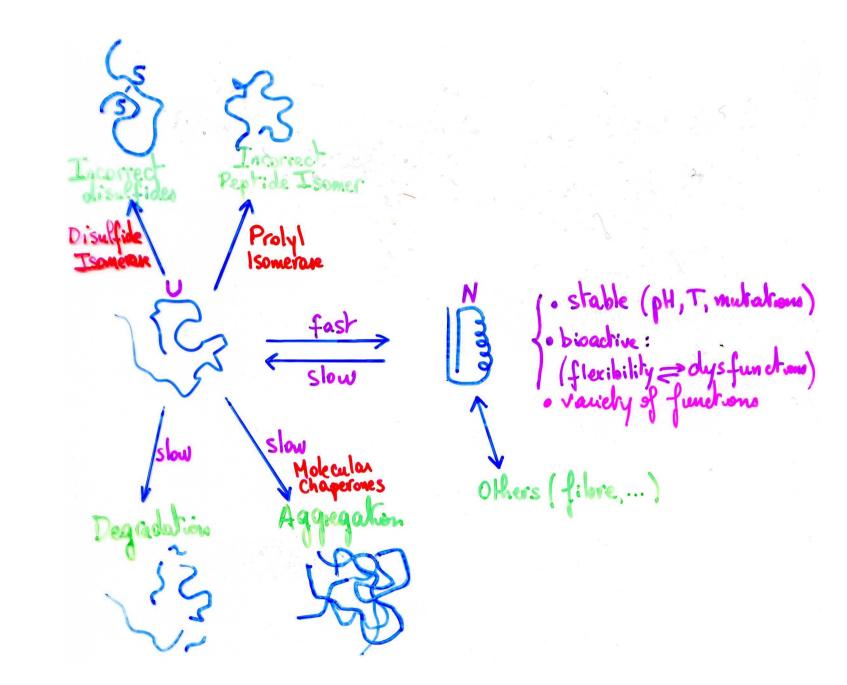
M1 Spécialité Bioinformatique Lecture 1A

P. Derreumaux

The protein folding problem: Generalities



The Protein Folding problem 2 aspects 5 structures from sequences. Sequence events from divordered
 to native structures (characterization
 of TSE) ->
 sequence elements coding for poteni
 topology

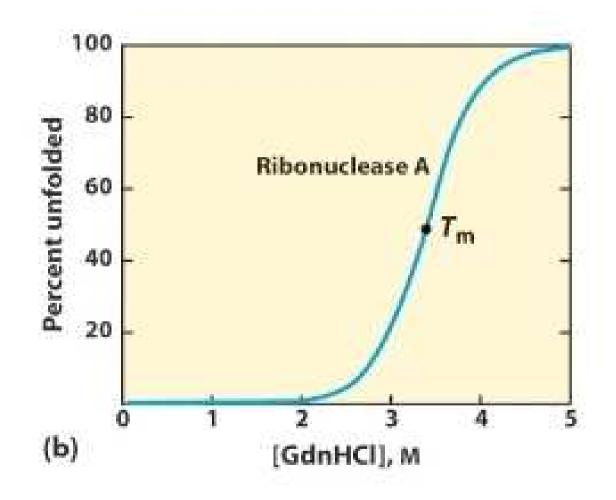


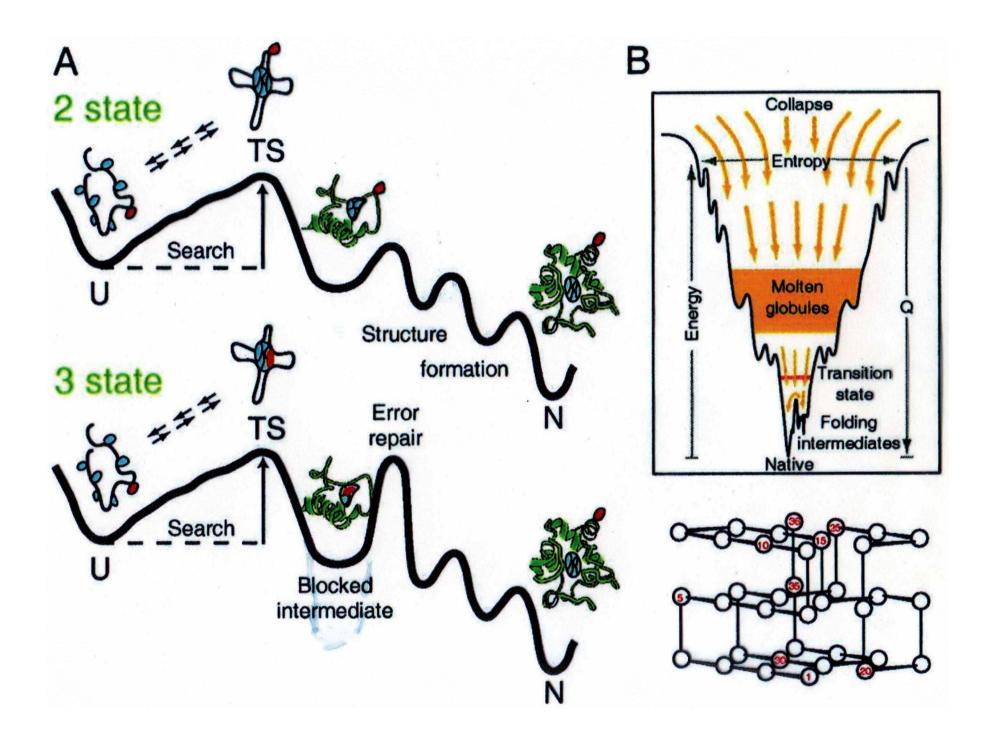
-Folding under chaperon control

-Thermodynamic hypothesis of folding (Anfinsen)

- Amyloid Fibril Formation

Monomeric Protein Folding: a simple experimental view





Monomeric Protein Folding: a simple view

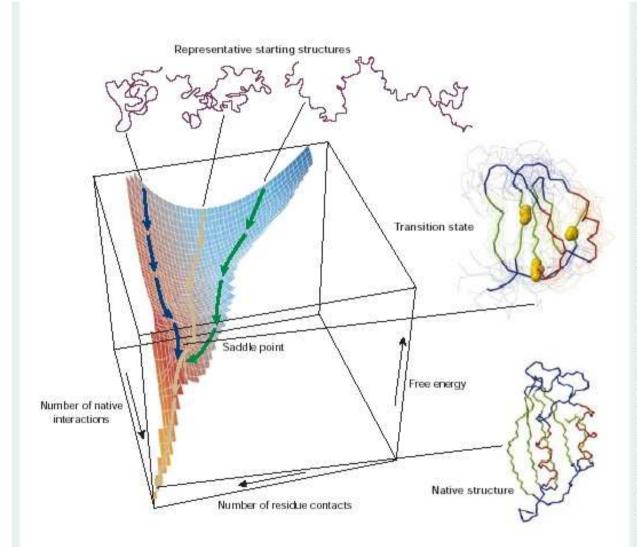
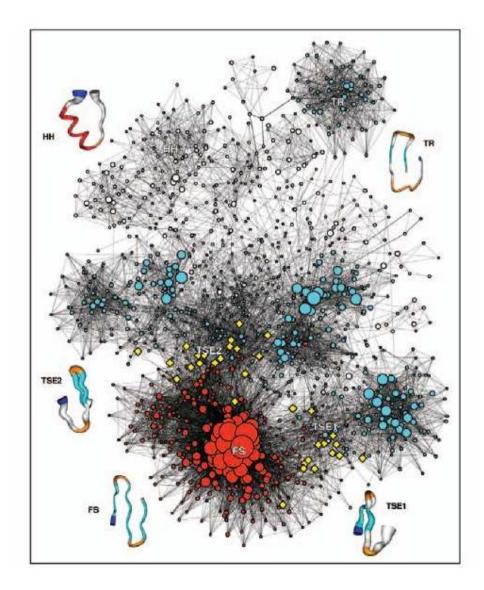
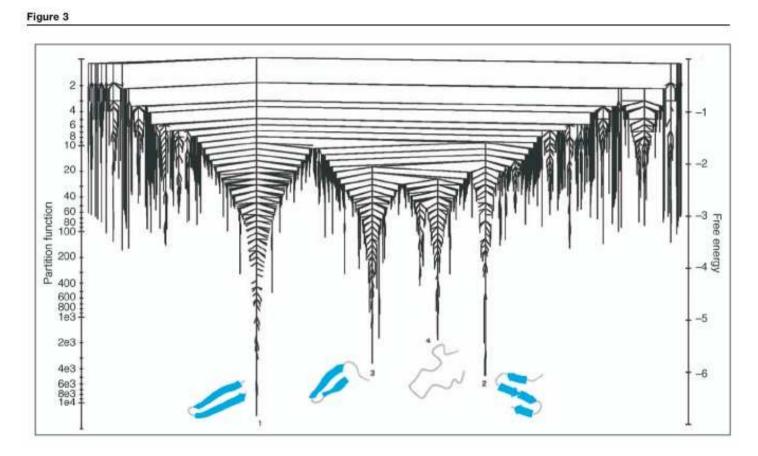


Figure 1 A schematic energy landscape for protein folding. The surface is derived from a computer simulation of the folding of a highly simplified model of a small protein. The surface 'funnels' the multitude of denatured conformations to the unique native structure. The critical region on a simple surface such as this one is the saddle point corresponding to the transition state, the barrier that all molecules must cross if they are to fold to the native state. Superimposed on this schematic surface are ensembles of structures corresponding to different stages of the folding process. The transition state ensemble was calculated by using computer simulations constrained by experimental data from mutational studies of acylphosphatase18. The yellow spheres in this ensemble represent the three 'key residues' in the structure; when these residues have formed their native-like contacts the overall topology of the native fold is established. The structure of the native state is shown at the bottom of the surface; at the top are indicated schematically some contributors to the distribution of unfolded species that represent the starting point for folding. Also indicated on the surface are highly simplified trajectories for the folding of individual molecules. Adapted from ref. 6.

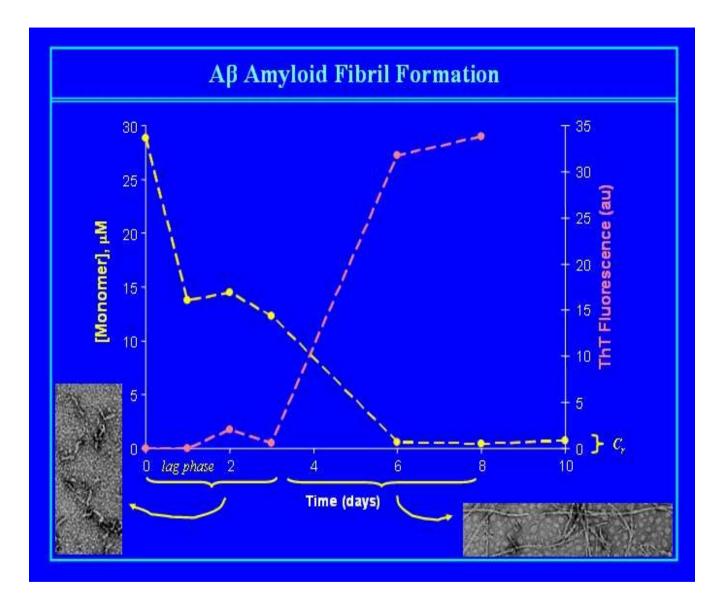
The funnel energy surface is oversimplified



The funnel energy surface is oversimplified

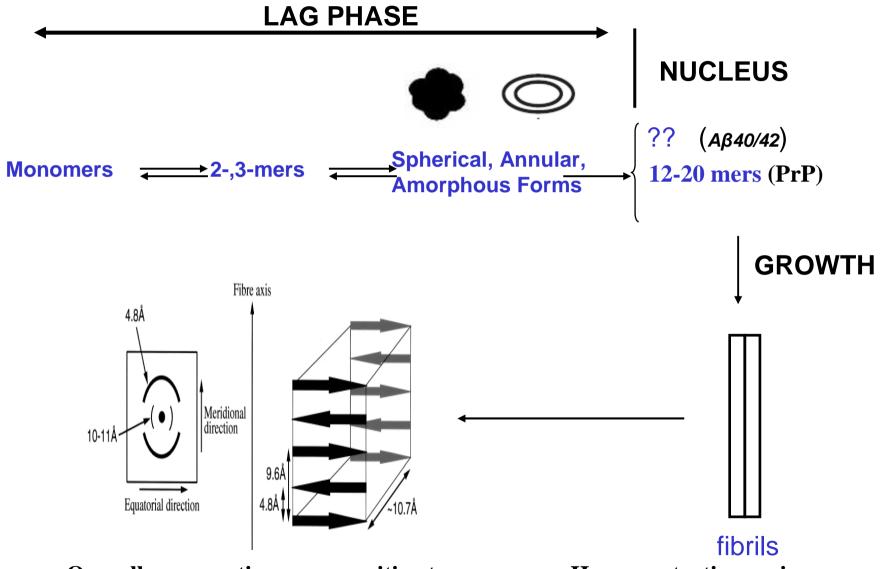


Transition disconnectivity graph of a β -hairpin (the C-terminal segment from the B1 domain of protein G). A total of 4 μ s implicit solvent molecular dynamics simulations at 360 K were sampled to obtain a sufficient number of folding-unfolding events [29]. Representative structures of the deepest free energy minima are shown and labeled 1–4. The left vertical axis shows the partition function of the minima and barriers. The right vertical axis shows the free energy of the minima and barriers. Reproduced with permission from [29].



Amyloid fibril formation is characterized by a polymerization-nucleation process

Details of the polymerization-nucleation process



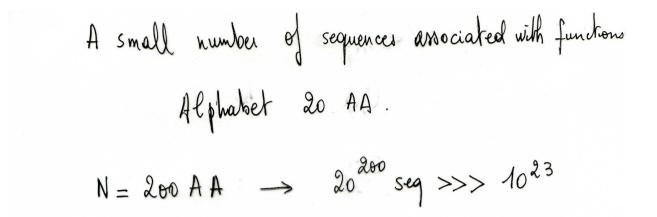
Overall, aggregation very sensitive to sequence, pH, concentration, anions

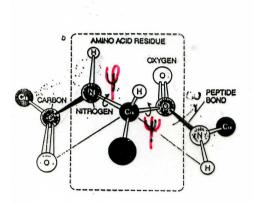
Prediction 3D structure from sequence

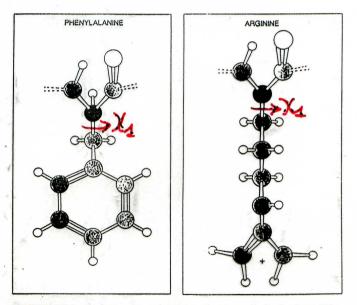
-Genomic Programs and the number of sequences

-Structures are more conserved than sequences

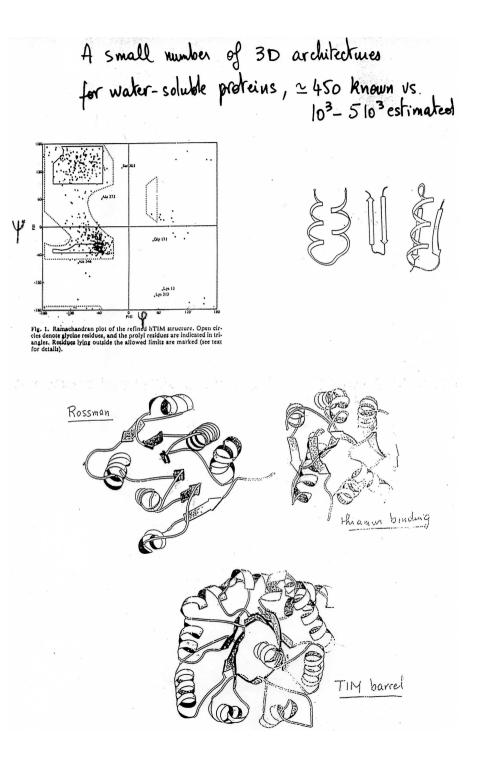
-Experimental costs for structure determination







DIFFERENCES in the shape, size and polarity of amino acids derive from differences in their side chains. In phenylalanine, for example, the side chain is nonpolar and cyclic, whereas the side chain of arginine is both strongly polar and linear.



domain residues 213 to 326; Ammon et al., 1900), (g) immunoglobulin, 3hhr (chain B, domain residues 32 to 131; De Vos et al., 1992); (h) UB, 1sha (chain A; Waksman et al., 1992); (i) jelly roll, 2stv (Fridborg et al., 1965; (j) plaitfold, 1aps (Saudek et al., 1989).

Table 1. Summary of average % of secondary structure residues which are in supersecondary structural units and motifs in all the categories of domain structure

		ββ, αα, βαβ	% Secondary structure residue +(βαββ)	s in: (β ₄)
Fold Globin UpDown Trefoil TIM barrel OB folds Doubly wound Ig UB roll Jelly roll Plaitfold All superfolds Non-superfolds	n 4 6 4 13 5 38 17 6 6 6 23 122 516	88.0 90.3 83.4 82.4 76.7 68.1 66.7 55:3 47.3 37.8 64.7 61.7	88.0 90.3 83.4 83.4 76.7 70.3 66.7 55.3 47.3 62.8 70.5 63.4	88.0 90.3 90.2 83.4 81.4 70.3 85.6 61.1 67.9 64.0 74.8 66.3

Non-superious β_{10} The *n* value is the number of domains per fold. The $\beta\beta$, $\alpha\alpha$, $\beta\alpha\beta$ column is the mean value for % of secondary structure residues in supersecondary in supersecondary structural units. The +($\beta\alpha\beta\beta$) column is the mean value for % of secondary structure residues in supersecondary structural units and in the $\beta\alpha$ -Greek key motif. The + β_4 column is the mean value for % of secondary structure residues in supersecondary structural units and in the $\beta\alpha$ -Greek key and in the β_4 Greek key motifs. This is the order in which the motif additions were made to the data.

PDB 1999 ~ 50% secondary structures

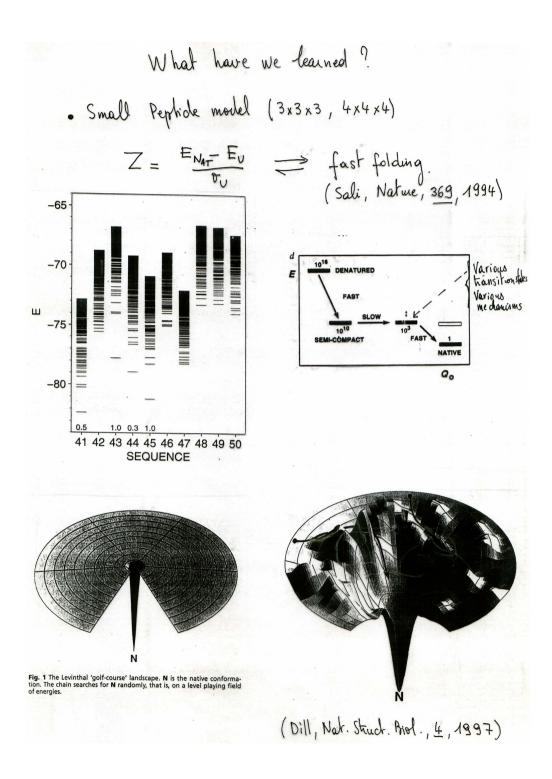
The Protein Folding problem and A limited number of folds. (~ 800< Nurston)

baric reasons: elsvious - structural stability $E_{Nat} \rightarrow P(E_{Nat}) = e^{-\frac{E_{Nat}}{kT}} \gg P(E_{i})$ $e^{-\frac{E_{Nat}}{kT}} \gg P(E_{i})$

IDP: Intrinsincally Disordered Proteins

- 30% of eukaryotic genome encoded proteins
- disordered over the full or part of the sequence
- fold upon binding on one or multiple partners

Thermodynamic, kinetic and dynamic ameto Mostrly lattice Models. (Cx) Off-lattice Ca Models (Nymeyer, DNAS, 35, 1998) Kinetic Properties. Thermalynamic Properties full enumeration NC, GEM known Canonical Eusenble run MC Hetropolis min (1, e-AE KBT) $Q = \sum_{i=1}^{N_{\rm C}} e^{-\frac{E_i}{k_{\rm T}}}$ moves . A = - KOT IN Q=E-TS $P_{Not} = \frac{e^{-\frac{E_{NaT}}{K_{0T}}}}{Q} \rightarrow T_{f}$ Figure 3. Crankshaft motion. The original is indicated by solid lines and filled circles. conformations are indicated by dashed line $\downarrow c_{v} \rightarrow T_{\theta} \quad \downarrow$ • $\mathcal{N}(4) = f(4)$ in function of sequence, T, PEF $\begin{array}{c} L_{3} S \longrightarrow T_{q} \\ (\frac{\partial S}{\partial T} = 0) \end{array}$ Nc = N (30-mers 2D, 15-mers 3D) = N_{compoct} (27-mers, 3D)



Coût d'une structure 3D

	Coût moyen	Change de succès
Protéines solubles bactériennes	\$140000	35%
Protéines humaines solubles (kinases, protéases,)	\$450000	35%
Protéines membranaires bactériennes	\$1,5 million	10%
Protéines membranaires humaines	\$2,5 million	10%

R.C. Stevens, Drug Discovery, 2003 Coût n'incluant pas les développements technologiques ni l'amortissement des équipements lourds M1 Spécialité Bioinformatique Lecture 1B

P. Derreumaux

Predicting Protein Structures from AA:

Two perspectives

- 1. If you were a theoretical chemist or physicist
- 2. If you were a pure bioinformatician

From Quantum Mechanics. $H\Psi = E\Psi$ to Molecular Mechanicas and Molecular Aynamics syst: ensemble of $\vec{F} = m\vec{a}$.

Niveau d'approximation en Mécanique Marique

Born-Oppenheimer : les mouvements des électrons sont découplés des mouvements des noyaux.

Les électrons sont représentés par leurs effets : charges partielles, paramètres (distance de référence entre 2 atomes liés de façon covalente, ...).

Seuls, les mouvements des noyaux sont considérés et sont traités dans un modèle de *mécanique classique*.

 \rightarrow champs de forces semi-empiriques

 \rightarrow grands systèmes, typiquement < 10000 atomes

 \rightarrow solutions numériques

 \rightarrow dynamique, limitée à ~ 1 ns (10⁻⁹ s)

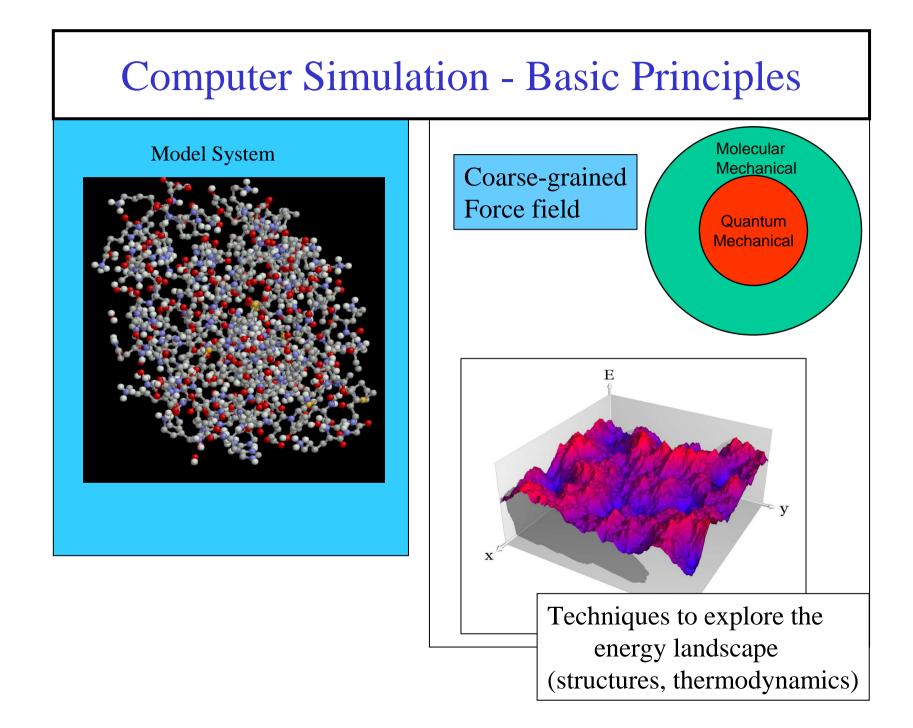
 \rightarrow représentation du solvant et des contre-ions

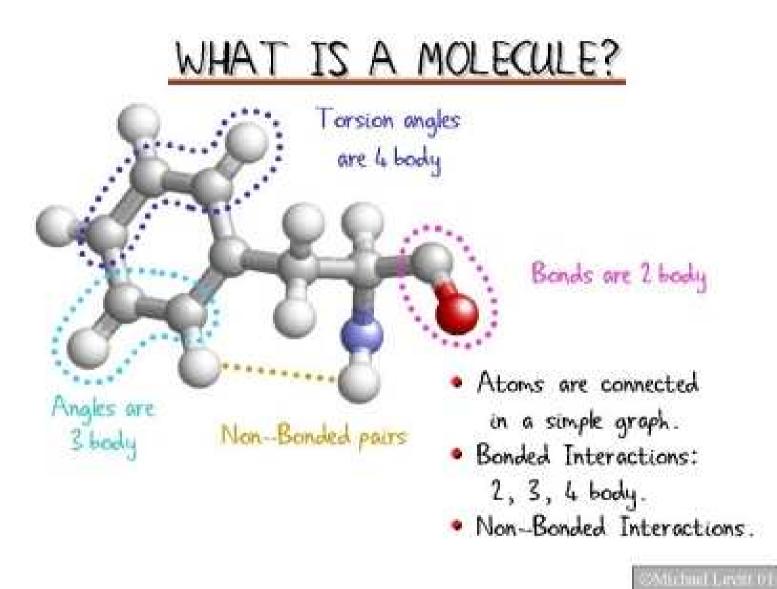
 \rightarrow énergie conformationnelle

 \rightarrow pas de réactions chimiques, polarisation, transfert d'électrons

δq

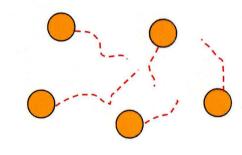
Capacité de transfert : des paramètres calculés et testés sur de petites molécules peuvent être utilisés pour des systèmes plus importants.



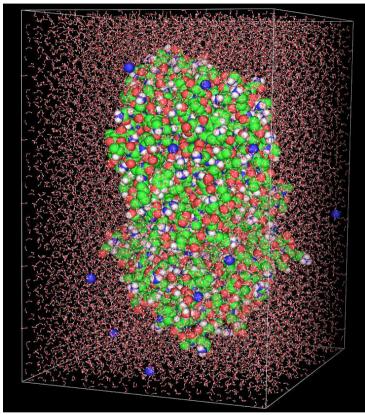


What is a molecular dynamics simulation?

- Simulation that shows how the atoms in the system move with time
- Typically on the nanosecond timescale
- Atoms are treated like hard balls, and their motions are described by Newton's laws.

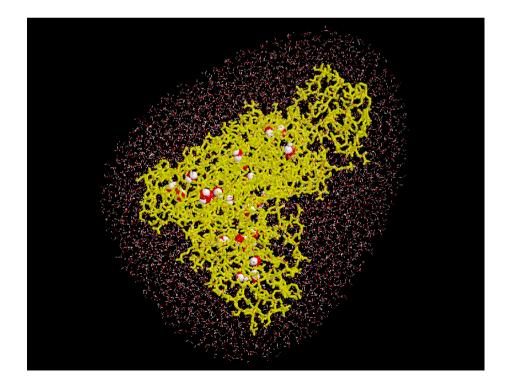


Treatment solvent

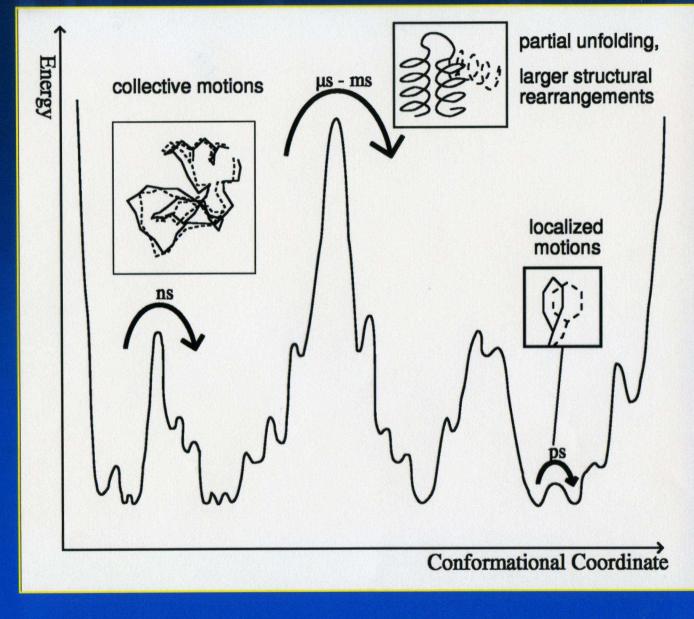


Box or Droplet ? SPC, TIP3P..?

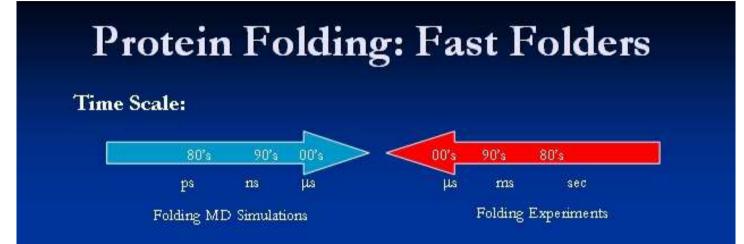
explicit or implicit?



Proteins jump between many, hierarchically ordered conformational substates



H. Frauenfelder et al., Science 229 (1985) 337



- Trp-cage, designed mini-protein (20 aa): 4μs
- β-hairpin of C-terminus of protein G (16 aa) : 6µs
- Engrailed homeodomain (En-HD) (61 aa): ~27μs
- WW domains (38-44 aa): >24μs
- **Ξ** Fe(II) cytochrome b₅₆₂ (106 aa): extrapolated ~5μs
- B domain of protein A (58 aa): extrapolated ~8μs

Ideally: MD simulations in solvent. $\begin{array}{l} \text{MD simulation}: a simple application.} \\ \text{F} = m a = m \frac{dv}{dt} = m \frac{dk}{dt^2}. \end{array}$ $supp a = dr = \frac{dv}{dt}$ $v = at + v_o$ $x = vt + x_0 = at^2 + v_0 t + z_0$ with $a = -\frac{1}{m} \frac{dV}{dx}$.

To calculate a trajectory, one needs. - mitrial pontions of the particles - mitrial distribution of velocities - the gradient of the patential energy function

Molecular Dynamics Simulation

Molecule: (classical) N-particle system

Newtonian equations of motion:

with
$$m_i \frac{d^2}{dt^2} \vec{r}_i = \vec{F}_i(\vec{r})$$
$$\vec{F}_i(\vec{r}) = -\nabla_i V(\vec{r})$$

Integrate numerically via the "leapfrog" scheme

$$\boldsymbol{v}(t + \frac{\Delta t}{2}) = \boldsymbol{v}(t - \frac{\Delta t}{2}) + \frac{\boldsymbol{F}(t)}{m}\Delta t$$
$$\boldsymbol{r}(t + \Delta t) = \boldsymbol{r}(t) + \boldsymbol{v}(t + \frac{\Delta t}{2})\Delta t$$

with ∆t ≈ 1fs!

(equivalent to the Verlet algorithm)

 $\Delta t = 1 - 2 fs$

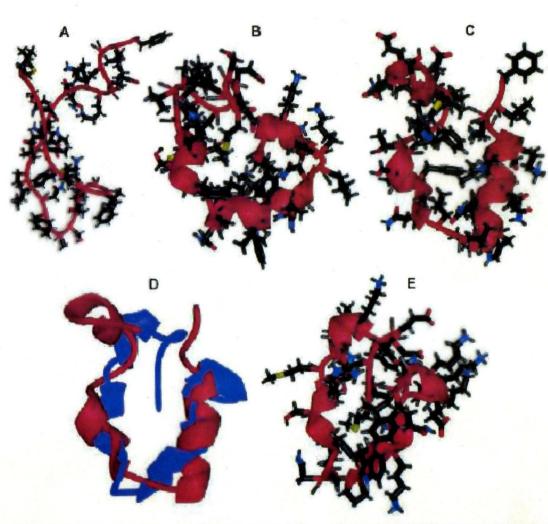
Molecular dynamics

36 amino acids +3,000 water molecules

1-ps simulation

256 parallel processors on a CRAY T3E

2 months of computer time



Y.Duan and P.A.Kollman, Science, 282,740 (1998)

Fig. 1. Ribbon representations of (A) the unfolded, (B) partially folded (at 980 ns), and (C) native structures, and (E) a representative structure of the most stable cluster and (D) the overlap of the native (red) and the most stable cluster (blue) structures, generated with UCSF MidasPlus. Color code [except (D)]: red, main chain atoms and oxygen; black, non-main chain carbon; blue, non-main chain nitrogen; gray, hydrogen; yellow, sulfur.

Atomic-Level Characterization of the Structural Dynamics of Proteins

David E. Shaw,^{1,2}* Paul Maragakis,¹† Kresten Lindorff-Larsen,¹† Stefano Piana,¹† Ron O. Dror,¹ Michael P. Eastwood,¹ Joseph A. Bank,¹ John M. Jumper,¹ John K. Salmon,¹ Yibing Shan,¹ Willy Wriggers¹

Molecular dynamics (MD) simulations are widely used to study protein motions at an atomic level of detail, but they have been limited to time scales shorter than those of many biologically critical conformational changes. We examined two fundamental processes in protein dynamics—protein folding and conformational change within the folded state—by means of extremely long all-atom MD simulations conducted on a special-purpose machine. Equilibrium simulations of a WW protein domain captured multiple folding and unfolding events that consistently follow a well-defined folding pathway; separate simulations of the protein's constituent substructures shed light on possible determinants of this pathway. A 1-millisecond simulation of the folded protein BPTI reveals a small number of structurally distinct conformational states whose reversible interconversion is slower than local relaxations within those states by a factor of more than 1000.

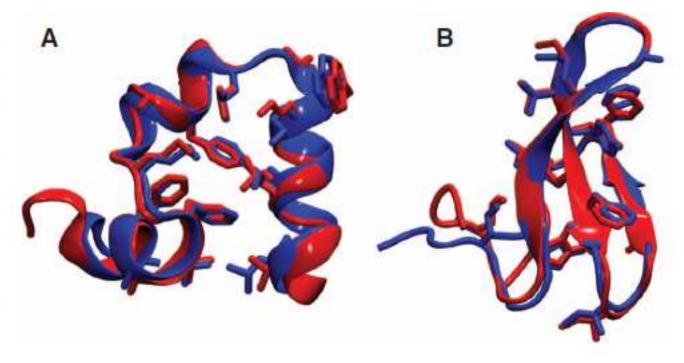
Many biological processes involve functionally important changes in the threedimensional structures of proteins. Conformational changes associated with protein folding (1), signal transduction (2), the catalytic cycles of enzymes (3), and the operation of molecular machines and motor proteins (4) often involve transitions among two or more structurally distinct sta terized as "ba "energy lands: Substantial both experime techniques, in c and the ways th them. It has pa ally characteriz states and to el anisms involve states.

All-atom n tions are desig view of the n cules (9), prod the potential t shots generated tational constr

¹D. E. Shaw Reses 10036, USA. ²Cen matics, Columbia I *14, whom corres David Shaw@DES †These authors or

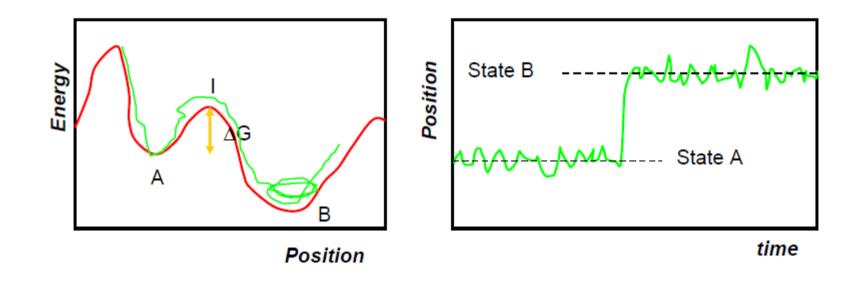
www.sciencemag.org SCIENCE VOL 330 15 OCTOBER 2010

Fig. 1. Folding proteins at x-ray resolution, showing comparison of x-ray structures (blue) (*15*, *24*) and last frame of MD simulation (red): (**A**) simulation of villin at 300 K, (**B**) simulation of FiP35 at 337 K. Simulations were initiated from completely extended structures. Villin and FiP35 folded to their native states after 68 μs and 38 μs, respectively,



and simulations were continued for an additional 20 µs after the folding event to verify the stability of the native fold.

Crossing energy barriers



The actual transition time from A to B is very quick (a few pico seconds).

What takes time is waiting. The average waiting time for going from A to B can be expressed as:

$$\tau_{A\to B} = C e^{\frac{\Delta G}{kT}}$$

1